

# ME-662 CONVECTIVE HEAT AND MASS TRANSFER

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LECTURE-38 CONV M T - COUETTE FLOW MODEL

# LECTURE-38 CONV M T - COUETTE FLOW MODEL

- 1 Gas Injection - Effect of property variation and  $\omega_T$  - LBL
- 2 Gas Injection - Effect of property variation and  $\omega_T$  - TBL
- 3 Benzene evaporation in convective environment

Couette flow model permits effects of fluid property variations to be studied.



## Soln ( Contd ) - 1 - L38( $\frac{2}{14}$ )

$$N_w = \frac{\rho_m D \partial \omega_g / \partial y|_w}{\omega_{g,w} - \omega_{g,T}}$$

Hence

$$N_w (\omega_{g,y} - \omega_{g,w}) = \rho_m D \frac{\partial \omega_g}{\partial y} \Big|_y - N_w (\omega_{g,w} - \omega_{g,T}) \quad \text{or}$$

$$\rho_m D \frac{\partial \omega_g}{\partial y} \Big|_y = N_w (\omega_{g,y} - \omega_{g,T})$$

where  $D = \text{const} \neq F(\omega_g)$  because  $p$  &  $T$  are const, but

$$\begin{aligned} \rho_m &= \frac{p}{R_u T} M_{mix} = \frac{p}{R_u T} \left( \sum \frac{\omega_j}{M_j} \right)^{-1} \\ &= \frac{p}{R_u T} \left[ \frac{M_g M_a}{M_a \omega_g + M_g (1 - \omega_g)} \right] \end{aligned}$$

## Soln ( Contd ) - 2 - L38( $\frac{3}{14}$ )

Substitution and integration from  $y = 0$  to  $\delta$  gives

$$\int_{\omega_{g,w}}^0 \frac{d\omega_g}{a\omega_g^2 + b\omega_g + c} = \frac{N_w R_u T \delta}{p M_g M_a D} \text{ with}$$

$$a = (M_a - M_g), \quad b = M_g - \omega_{g,T} (M_a - M_g), \quad c = -M_g \omega_{g,T}$$

where the LHS is given by

$$\begin{aligned} \text{LHS} &= \frac{1}{\sqrt{b^2 - 4ac}} \ln \left[ \frac{2a\omega_g + b - \sqrt{b^2 - 4ac}}{2a\omega_g + b + \sqrt{b^2 - 4ac}} \right]_{\omega_{g,w}}^0 \\ &= \frac{1}{M_g + \omega_{g,T} (M_a - M_g)} \ln \left[ 1 + B + \omega_{g,T} B \left( \frac{M_a}{M_g} - 1 \right) \right] \end{aligned}$$

where

$$B = \frac{0 - \omega_{g,w}}{\omega_{g,w} - \omega_{g,T}} = \frac{\omega_{g,w}}{\omega_{g,T} - \omega_{g,w}} \quad \text{and} \quad \omega_{g,w} = \omega_{g,T} \times \frac{B}{1 + B}$$

## Soln ( Contd ) - 3 - L38( $\frac{4}{14}$ )

Now, for the Couette flow model

$$N_w = g B, \quad \text{and} \quad \frac{R_u T}{\rho M_g} = \frac{1}{\rho_g}$$

Therefore

$$\text{RHS} = \frac{N_w R_u T \delta}{\rho M_g M_a D} = \frac{g B \delta}{\rho_g M_a D}$$

Equating LHS = RHS and rearranging

$$\left( \frac{g \delta}{\rho_g D} \right) = \left( \frac{M_a}{M_g} \right) \left[ \frac{\ln(1 + B^*)}{B^*} \right] \quad \text{where}$$

$$B^* = B \left\{ 1 + \omega_{g,T} \left( \frac{M_a}{M_g} - 1 \right) \right\}. \quad \text{Hence}$$

$$\left( \frac{g}{g^*} \right)_{vp} = \frac{\ln(1 + B^*)}{B^*} \quad (\text{Ans}) \rightarrow \left( \frac{g}{g^*} \right)_{cp} = \frac{\ln(1 + B)}{B}$$

where subscript 'vp' for variable and 'cp' for const property.

# Soln - $\left(\frac{g}{g^*}\right) \sim B$ for $\omega_{g,T} = 1$ - L38 $\left(\frac{5}{14}\right)$

B	cp	$vp_{CO_2}$	$vp_{He}$	$vp_{H_2}$	$\omega_{g,w}$
0	1.0	1.0	1.0	1.0	0.0
.25	.893	.926	.571	.422	.200
.50	.811	.864	.422	.291	.333
1.0	.693	.768	.291	.189	.500
1.5	.611	.695	.228	.144	.600
2.0	.549	.638	.189	.117	.667
2.5	.501	.591	.163	.0998	.714
3.0	.462	.552	.144	.0873	.750

- 1  $\omega_{g,T} = 1$  implies that the gas is the only transferred substance . Also,  $B^* = B M_a / M_g$ .
- 2  $(g/g^*)_{vp,CO_2} > (g/g^*)_{cp}$  because  $M_{CO_2} > M_{air}$
- 3 For He and  $H_2$ , this trend reverses.
- 4  $\omega_{g,w}$  increases with B

# Soln - $\left(\frac{g}{g^*}\right) \sim B$ for $\omega_{g,T} = 0.01$ - L38( $\frac{6}{14}$ )

B	cp	$vp_{CO_2}$	$vp_{He}$	$vp_{H_2}$	$\omega_{g,w}$
0	1.0	1.0	1.0	1.0	0.0
.25	.893	.893	.887	.888	.002
.50	.811	.811	.802	.792	.0033
1.0	.693	.694	.681	.668	.005
1.5	.611	.612	.598	.584	.006
2.0	.549	.550	.536	.522	.0067
2.5	.501	.502	.488	.474	.0071
3.0	.462	.463	.449	.435	.0075

- 1  $\omega_{g,T} = .01$  implies that the gas in the transferred substance is a small fraction - rest is air.
- 2  $(g/g^*)_{vp,CO_2} \simeq (g/g^*)_{cp}$
- 3 For He and  $H_2$ ,  $(g/g^*)_{vp} < (g/g^*)_{cp}$
- 4  $\omega_{g,w}$ , though small, increases with B



# Correlation with $\left(\frac{M_{mix,\infty}}{M_{mix,w}}\right)$ - L38( $\frac{7}{14}$ )

Here,  $M_{mix,w} = M_a M_g / (M_a \omega_{g,w} + M_g (1 - \omega_{g,w}))$   
and  $M_{mix,\infty} = M_a$  ( because  $\omega_{g,\infty} = 0$  ). Hence, from slide 4,  
and using  $\omega_{g,w} = \omega_{g,T} \times B / (1 + B)$

$$B^* = B \left\{ 1 + \omega_{g,T} \left( \frac{M_a}{M_g} - 1 \right) \right\}.$$

$$\frac{B^*}{B} = 1 + \left( \frac{1 + B}{B} \right) \left( \frac{M_{mix,\infty}}{M_{mix,w}} - 1 \right)$$

$$\frac{(g/g^*)_{vp}}{(g/g^*)_{cp}} = \frac{\ln(1 + B^*)}{B^*} \times \frac{B}{\ln(1 + B)}$$

This shows dependence on  $M_{mix,w}/M_{mix,\infty}$  and B as recommended correction from boundary layer flow model.  
If  $\omega_{g,T} = 0$ ,  $B^* = B$ . If  $\omega_{g,T} = 1$ ,  $B^* = B (M_a/M_g)$

# Turbulent Couette Flow - 1 - L38( $\frac{8}{14}$ )

Here, the governing Eqn will be

$$N_w (\omega_g - \omega_{g,T}) = \rho_m (D + D_t) \frac{d\omega_g}{dy}$$

where

$$\rho_m D_t = \rho_m \frac{\nu_{t,ref}}{Sc_t} \quad \text{But, from Van-Driest model}$$

$$\nu_{t,ref} = \frac{\mu_t}{\rho_{ref}} = l_m^2 \frac{\partial u}{\partial y} \rightarrow \frac{\partial u}{\partial y} = C$$

$$= C \left( \frac{\nu_{ref}}{u_\tau} \right)^2 (\kappa y^+)^2 \left\{ 1 - \exp\left(-\frac{y^+}{A^+}\right) \right\}^2 \quad \text{and}$$

$$= C \left( \frac{\nu_{ref}}{u_\tau} \right)^2 (0.08 \delta^+)^2 \quad \text{for } y^+ > 26 \quad \text{where}$$

$$C \left( \frac{\nu_{ref}}{u_\tau} \right)^2 = C \frac{\nu_{ref}^2 \rho_{ref}}{\tau_w} = C \times \frac{\mu_{ref} \nu_{ref}}{\mu_{ref} C} = \nu_{ref}$$

# Turbulent Couette Flow - 2 - L38( $\frac{9}{14}$ )

Substituting for  $D_t$  and  $\rho_m$ , we have

$$\begin{aligned} N_w (\omega_g - \omega_{g,T}) &= \rho_m D \left(1 + \frac{\nu_{t,ref}}{Sc_t D}\right) \frac{d\omega_g}{dy} \\ &= \left(\frac{D p M_a M_g}{R_u T}\right) \times \frac{u_\tau / \nu_{ref}}{M_a \omega_g + M_g (1 - \omega_g)} \\ &\times F \times \frac{d\omega_g}{dy^+} \quad \text{where} \end{aligned}$$

$$\begin{aligned} F &= 1 + \left(\frac{Sc}{Sc_t}\right) (\kappa y^+)^2 \left\{1 - \exp\left(-\frac{y^+}{A^+}\right)\right\}^2 \quad y^+ < 26 \\ &= 1 + \left(\frac{Sc}{Sc_t}\right) (0.08 \delta^+)^2 \quad y^+ > 26 \end{aligned}$$

# Turbulent Couette Flow - 3 - L38( $\frac{10}{14}$ )

Taking  $N_w = g B$ ,  $(\rho M_g)/(R_u T) = \rho_g$  and  $u_\tau = U_\infty \sqrt{C_{f,x}/2}$ ,

$$\text{LHS} = \left( \frac{g}{\rho_g U_\infty} \sqrt{\frac{2}{C_{f,x}}} Sc \right) \times \text{INT where INT} = \int_0^{\delta^+} \frac{dy^+}{F}$$

$$\begin{aligned} \text{RHS} &= \frac{M_a}{B} \int_{\omega_{g,w}}^0 \frac{d\omega_g}{(\omega_g - \omega_{g,T}) \{M_a \omega_g + M_g (1 - \omega_g)\}} \\ &= \frac{\ln(1 + B^*)}{B^*} \rightarrow B^* = B \left\{ 1 + \omega_{g,T} \left( \frac{M_a}{M_g} - 1 \right) \right\} \end{aligned}$$

Taking  $A^+ = 26$  and  $Sc_t = 0.9$ , we have

INT = 9.62 for  $CO_2 - Air$ ,  $Sc = 0.96$

INT = 14.57 for  $H_2 - Air$  and He-Air,  $Sc = 0.22$

# Turbulent Couette Flow - 4 - L38( $\frac{11}{14}$ )

Therefore

$$\frac{g_{vp}}{\rho_g U_\infty} \times \sqrt{\frac{2}{C_{f,x}}} \times Sc = \frac{1}{INT} \times \frac{\ln(1 + B^*)}{B^*}$$

and

$$\frac{(g/g^*)_{vp}}{(g/g^*)_{cp}} = \frac{\ln(1 + B^*)}{B^*} \times \frac{B}{\ln(1 + B)}$$

This result is same as that for a Laminar boundary layer. This is because it is assumed that the value of INT is same for 'cp' and 'vp' conditions.

Note that  $g_{vp}$  is significantly influenced by INT ( Sc ).

## Evaporation of $C_6H_6$ - L38( $\frac{12}{14}$ )

**Prob:**  $C_6H_6$  evaporates from the outer surface of a circular cylinder in air flowing at 6 m/s normal to the cylinder.

From expts,  $h_{cof, v_w=0} = 85 \text{ W/m}^2\text{-K}$  and  $B = 0.9$ .

Allowing for property variations, estimate  $N_w$  and  $\omega_w$ .

Given:  $Sc = 1.71$ ,  $Pr = 0.71$ ,  $cp_{C_6H_6} = 1.69 \text{ kJ/kg-K}$  and  $cp_a = 1.01 \text{ kJ/kg-K}$ .

**Soln:** Here,

$$B = \frac{\omega_{v,\infty} - \omega_{v,w}}{\omega_{v,w} - 1} = 0.9 \rightarrow \omega_{v,w} = 0.4737 \quad (\text{Ans})$$

Therefore,  $\omega_{v,m} = 0.5 (\omega_{v,\infty} + \omega_{v,w}) = 0.2368$ .

$c_{pm} = 1.69 \times 0.2368 + 1.01 \times 0.7632 = 1.171 \text{ kJ/kg-K}$ .

Hence,  $g^* = (h_{cof, v_w=0} / c_{pm}) = 0.0726 \text{ kg/m}^2\text{-s}$ .

Also,  $M_{mix,\infty} = 29$  and

$M_{mix,w} = (0.4737/78 + 0.5263/29)^{-1} = 41.286$ .

## Soln ( Contd. ) - L38( $\frac{13}{14}$ )

For Flow over a cylinder<sup>1</sup>,  $Nu_{cp} \propto Pr^{0.37}$ .

Therefore, using the short-cut empirical formula

$$\begin{aligned}\frac{g_{vp}}{g_{cp}^*} &= \frac{\ln(1+B)}{B} \times \left(\frac{Pr}{Sc}\right)^{0.37} \times \left(\frac{M_{mix,\infty}}{M_{mix,w}}\right)^{-0.67} \\ &= \frac{\ln(1+0.9)}{0.9} \times \left(\frac{0.71}{1.71}\right)^{0.37} \times \left(\frac{29}{41.286}\right)^{-0.67} = 0.6525\end{aligned}$$

Therefore,  $g = 0.0726 \times 0.6525 = 0.0474 \text{ kg/m}^2\text{-s}$  ( Ans ) .

Thus, the effect of property variations is to reduce  $g_{vp}$  compared to  $g_{cp}$ .

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<sup>1</sup>Zhukauskas A *Heat Transfer from Tubes in Crossflow*,  
Eds: Hartnett J P and Irvine T F, Adv H T, vol 8, Academic Press, ( 1972 )

## Soln ( Contd. ) - L38( $\frac{14}{14}$ )

If we followed the Couette flow theory, then in this case,

$$B^* = B \left\{ 1 + \omega_{g,T} \left( \frac{M_a}{M_g} - 1 \right) \right\} = 0.3346$$

Hence

$$\left( \frac{g}{g^*} \right)_{vp} = \frac{\ln(1 + 0.3346)}{0.3346} = 0.8626$$

But, for variable properties,  $h_{cof, vp} = h_{cof, cp} \times Pr^{.25}$ .

Therefore,  $g_{vp} = g_{cp}^* \times (0.71)^{0.25} \times 0.8626 = 0.0575 \text{ kg/m}^2\text{-s}$ .

This value is **greater than** that obtained from the empirical formula. Thus, Couette flow theory provides an approximate answer due to linear velocity profile assumption.